

RESEARCH ARTICLE

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Key Points:

- Coseismic fault geometry and location closely match a preexisting blind fault
- Quaternary fold grows by distributed deformation in the overlying sediments and does not align with coseismic or early postseismic uplift
- Rheology contrast between the Tarim Basin and Tibet may explain the earthquake's low dip angle

Supporting Information:

- Supporting Information S1

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Blind Thrusting, Surface Folding, and the Development of Geological Structure in the M_w 6.3 2015 Pishan (China) Earthquake

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Abstract The relationship between individual earthquakes and the longer-term growth of topography and of geological structures is not fully understood, but is key to our ability to make use of topographic and geological data sets in the contexts of seismic hazard and wider-scale tectonics. Here we investigate those relationships at an active fold-and-thrust belt in the southwest Tarim Basin, Central Asia. We use seismic waveforms and interferometric synthetic aperture radar (InSAR) to determine the fault parameters and slip distribution of the 2015 M_w 6.3 Pishan earthquake—a blind, reverse-faulting event dipping toward the Tibetan Plateau. Our earthquake mechanism and location correspond closely to a fault mapped independently by seismic reflection, indicating that the earthquake was on a preexisting ramp fault over a depth range of ~9–13 km. However, the geometry of folding in the overlying fluvial terraces cannot be fully explained by repeated coseismic slip in events such as the 2015 earthquake nor by the early postseismic motion shown in our interferograms; a key role in growth of the topography must be played by other mechanisms. The earthquake occurred at the Tarim-Tibet boundary, with the unusually low dip of 21°. We use our source models from Pishan and a 2012 event to argue that the Tarim Basin crust deforms only by brittle failure on faults whose effective coefficient of friction is $\leq 0.05 \pm 0.025$. In contrast, most of the Tibetan crust undergoes ductile deformation, with a viscosity of order 10^{20} – 10^{22} Pa s. This contrast in rheologies provides an explanation for the low dip of the earthquake fault plane.

1. Introduction

The relationship between individual earthquake cycles and the long-term development of fold and thrust structures is not fully understood. Theoretical geometrical models have been proposed to describe how folds form and grow over time above thrust faults with various geometries (e.g., Erslev, 1991; Homza & Wallace, 1995; Suppe, 1983; Suppe & Medwedeff, 1990), and the development of thrust-related folds in nature over tens of thousands to millions of years has been studied by analysis of drainage, topography and structural geology (e.g., Burbank et al., 1999; Daëron et al., 2007). However, it is not clear how complex fault and fold structures develop through repeated earthquake cycles nor how these longer-term structures may or may not correlate with the surface deformation measured today (e.g., Nissen et al., 2007). The relative importance of interseismic, coseismic, and postseismic deformation to the development of geological structures and topography in regions of continental shortening is not known. Given that historical records of earthquakes are often much shorter than recurrence times for each fault, data sources with longer temporal coverage and global spatial coverage, such as digital elevation models and geomorphology, combined with dating, provide a valuable opportunity to augment assessments of seismic hazard (Stein & King, 1984; Taylor & Yin, 2009; Zhou et al., 2015; Talebian et al., 2016). Making use of this opportunity depends on developing an ability to interpret the spatial and temporal relationship between long-term structures and individual earthquakes. To address this issue, here we investigate structural growth at a blind range front thrust where we can directly compare surface motion observations from interferometric synthetic aperture radar (InSAR), subsurface seismic reflection data and surface geomorphology.

In addition, the Pishan earthquake represents an opportunity to investigate the wider question of what controls the geometry of active fault planes. By examining the material properties and behavior of the lithosphere

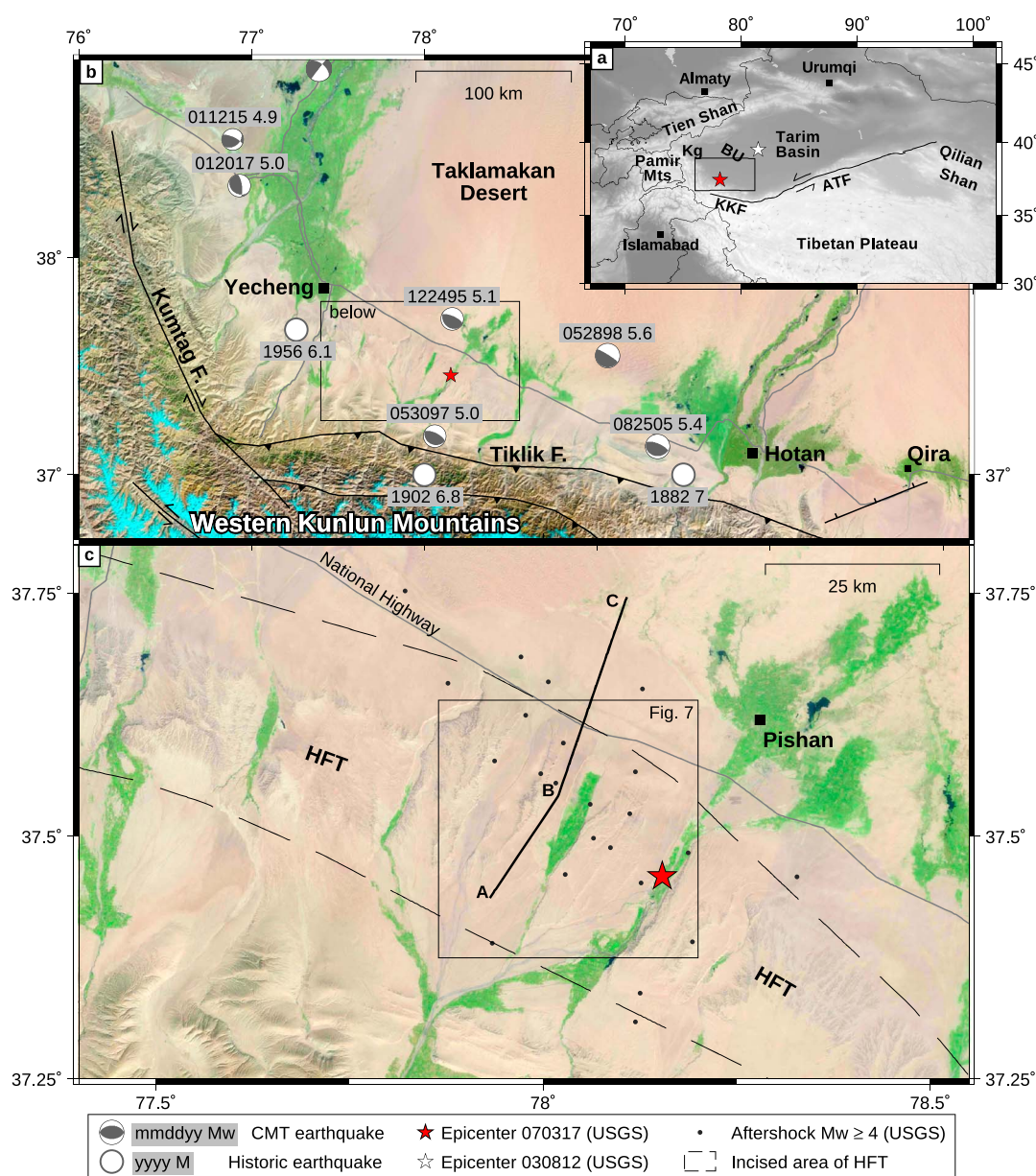


Figure 1. (a) Overview map of the Tarim Basin and its surroundings. ATF: Altyn Tagh Fault. KKF: Karakax Fault. BU: Bachu Uplift. Kg: Kashgar city. Black rectangle shows the area of Figure 1b. (b) Landsat 8 imagery (RGB 654) of the southwest Tarim Basin region. Faults after Avouac and Peltzer (1993) and Wei et al. (2013). (c) Landsat imagery around the epicenter of the 2015 Pishan earthquake (red star) and showing the incised area above the Hotan Fold and Thrust Belt (HFT). A-B-C shows the location of the seismic reflection profile (section 4).

beneath the Tarim Basin and Tibetan Plateau we are able to examine how the contrast between them affects the geometry of the resulting faulting.

On 3 July 2015 an M_w 6.3 earthquake hit Pishan county in Xinjiang, China, at 01:07:47 UTC (09:07 local time). Local media reported three deaths and damage to several thousand buildings. The epicenter was located approximately 20 km southwest of the historic Silk Road town of Pishan (also known as Guma), an oasis town on the edge of the Taklamakan Desert. Thirty aftershocks with magnitude $m_b \geq 4.0$ were reported by the United States Geological Survey (USGS) in the 7 months following the earthquake and form a distributed pattern roughly trending northwest-southeast, approximately 100 km in extent (Figure 1).

We present a source model for the Pishan earthquake based on Sentinel-1 interferometric synthetic aperture radar (InSAR) and teleseismic body wave modeling. We subsequently identify postseismic motion in

interferograms covering the first 7 months following the event. We correlate the relative locations and magnitudes of coseismic and postseismic displacements to the tectonic geomorphology of the area in order to constrain the mechanism by which the current topography has formed over the late Quaternary. We compare our results to a seismic reflection survey across the same location to investigate the full history of the underlying folds and thrusts in the Pliocene. Placing our results in a wider regional context, we use our source models for the Pishan earthquake and a smaller 2012 event 350 km to the northeast in the central Tarim Basin to constrain the rheology of components in the surrounding Tarim-Tibet collision zone and examine the effect they had on the geometry of the faulting at Pishan.

2. Tectonic Setting

Pishan is situated at the margin between the Tibetan Plateau and the southwestern corner of the Tarim Basin, 800 km north of the Himalayan front (Figure 1). The Tarim Basin acts as a rigid block within the India-Eurasia collision; although it is surrounded by mountain ranges undergoing active shortening—the Tien Shan to the north, the Pamir mountains to the west, and the Tibetan Plateau to the south—the Tarim Basin has low seismicity, flat topography, and GPS measurements show little internal deformation (England & Molnar, 2005; Molnar & Tapponnier, 1978). At the southwest corner of the Tarim Basin, where we focus our attention in this paper, gravity anomalies (Lyon-Caen & Molnar, 1984), receiver functions (Kao et al., 2001; Wittlinger et al., 2004), seismic refraction (Qiusheng et al., 2002), and geological mapping (Matte et al., 1996) show that the Tarim Basin is underthrusting the Western Kunlun Mountains, which form the edge of the Tibetan Plateau. GPS stations are relatively sparse, but 0–3 mm/yr of N-S convergence is possible within the uncertainties of the measurements (Ge et al., 2015) and the region has hosted M_s 7.0 and 6.8 earthquakes in 1882 and 1902, respectively, and an M_w 6.1 event in 1956 (Avouac & Peltzer, 1993; Lee et al., 1978; Li et al., 2012; United States Geological Survey, 2017) (Figure 1).

Across the northern range front of the Western Kunlun Mountains the elevation changes rapidly from 5,000 m to 2,100 m within a horizontal distance of just 50 km, and post-Mesozoic basin sediments have a thickness of around 10 km (Jiang & Li, 2014; Wei et al., 2013). The most basinward active shortening visible at the surface is in the low-lying, gently northward sloping piedmont. The structures here are known as the Hotan Fold and Thrust Belt and are expressed around Pishan as a belt of exposed and folded bedrock and terraces crossing from 77.0°E, 37.75°N to 78.75°E, 37.2°N (Figure 1). The total upper crustal shortening across the Hotan Fold and Thrust Belt has been of the order of several tens of kilometers (Jiang et al., 2013) and occurred in two stages, the first ~23–10 Ma and the present stage, into which the 2015 Pishan earthquake falls and which began ~5 Ma (Jiang & Li, 2014; Wei et al., 2013).

3. The 3 July Pishan Earthquake Fault Geometry and Coseismic Slip Distribution

3.1. Body Waveform Modeling

We performed a joint inversion of teleseismic P and SH waveforms to obtain the focal parameters and source time function of the Pishan mainshock. Seismograms (25 for P , 23 for SH) from stations at epicentral distances of 30–80° were band-pass filtered to 15–100 s and converted to the response of a WWSSN long-period seismometer. These were then inverted using the MT5 program of Zwick et al. (1994) (based on the algorithm of McCaffrey & Abers, 1988; McCaffrey et al., 1991). The method is well established and described elsewhere (e.g., Craig et al., 2011; Molnar & Lyon-Caen, 1989; Sloan et al., 2011). Stations were weighted to account for the uneven distribution of their azimuths, and P seismograms were given double the weight of SH seismograms to address the amplitude imbalance. The velocity structure at the source was specified as a simple half-space velocity model, with $V_p = 6.00$ km/s, $V_s = 3.45$ km/s. While this is likely to be fast relative to the uppermost crustal layers, it represents a good average for the bulk crustal properties, and at the frequency range used, the waveforms are not sensitive to short-length scale variations in the shallow velocity structure.

The solution with the best least squares fit to the data is shown in Table 1 and Figure 2. The stations have good azimuthal coverage, and the synthetic waveforms fit the observed seismograms well (Figure 2). Our result agrees with the focal mechanism in the GCMT catalog (Table 1) and shows a single reverse faulting earthquake on a fault striking WNW-ESE, parallel to the distribution of aftershocks. The waveform modeling cannot distinguish between the primary and auxiliary nodal planes but shows that the fault dips either steeply to the north or shallowly to the south (Figure 2). The centroid depth of the lowest misfit solution is 16.5 km, and the acceptable range of depths, based upon the percentage changes to the misfit, is ~10–19 km (Figure 2).

Table 1
Pishan Mainshock Fault Parameters From Catalogs, Body Wave Modeling, and Uniform Slip Inversion of InSAR

Model	Long. (deg)	Lat. (deg)	Strike (deg)	Dip (deg)	Rake (deg)	Centroid depth (km)	Length (km)	Width (km)	Slip (m)	M_o (10^{18} Nm)	M_w
GCMT	78.14 ^b	37.58 ^b	109	22	85	15.6	-	-	-	5.3	6.4
Auxiliary plane			294	68	92						
USGS NEIC	78.154 ^c	37.459 ^c	105	24	60	15.5	-	-	-	5.3	6.4
Auxiliary plane			317	69	102						
He et al. (2016) ^a	78.1390 ^b	37.7157 ^b	113.8	27.2	93.7	10.98	22.18	8	0.6	4.12	6.33
	± 0.0004	± 0.0005	± 0.2	± 0.2	± 0.6	± 0.05	± 0.05				
Wen et al. (2016) ^a	78.057 ^d	37.571 ^d	114.0	23.6	92.6	8.8	22.1	10.1	0.59	-	6.4
	± 0.4 km	± 0.3 km	± 1.6	± 1.5	± 3.2	± 0.4	± 0.5	± 1.0	± 0.06		
Body wave (this study)	78.14 ^e	37.58 ^e	102	23	73	16.5	-	-	-	2.5	6.2
Auxiliary plane			300	68	97						
InSAR (this study) ^a	78.140 ^f	37.777 ^f	112	21	83	11.0	21.5	9.3	0.6	4.0	6.3
	± 1.7 km	± 2.7 km	± 2.4	± 2.5	± 4.9	± 0.8	± 1.1	± 2.1	± 0.2	± 0.2	

^aUniform slip model. ^bCentroid. ^cEpicenter. Focal depth 20.0 km. ^dNot specified. ^eFixed at GCMT location. ^fUpdip projection of center.

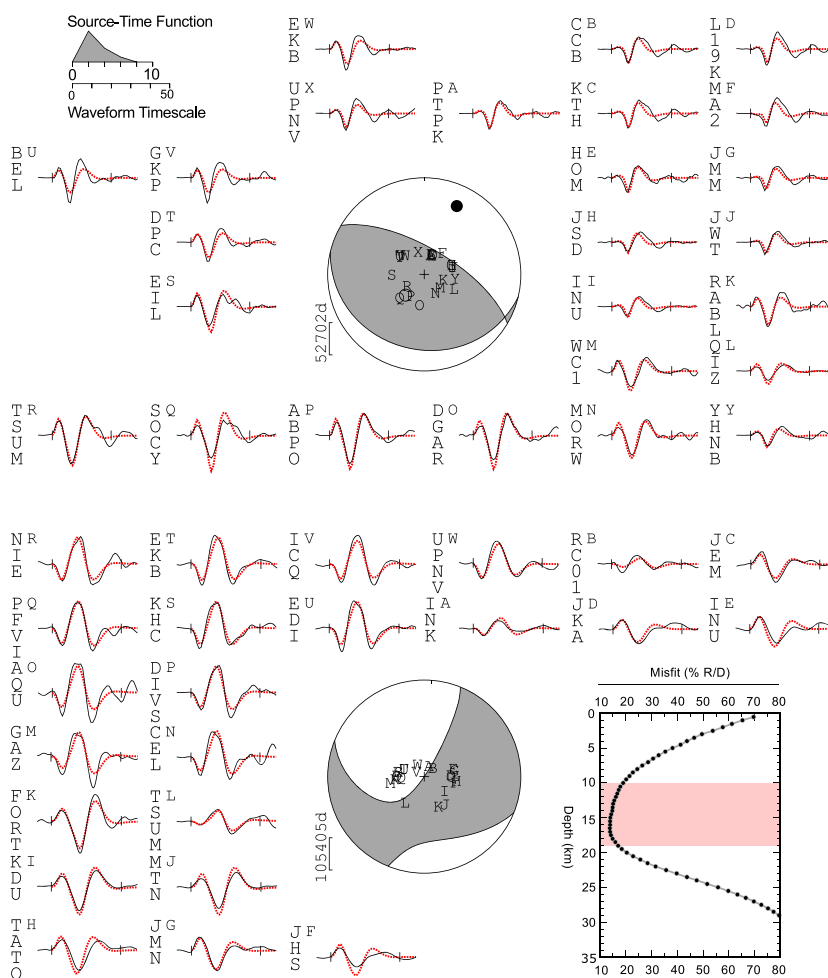


Figure 2. Focal mechanisms and source time function obtained from teleseismic body waveform modeling of (top) P and (bottom) SH waves. Mechanisms are shown as lower hemisphere projections with compressional quadrants shaded in gray. Seismograms are shown in black next to their station code, synthetic seismograms from the model solution are shown as dashed red lines. (bottom right) Misfits for solutions with depth fixed at a range of values, plotted as a percentage fraction of residual divided by data (% R/D). Pink band shows the acceptable range of depths referred to in section 3.1. Numerical values for this solution are listed in Table 1.

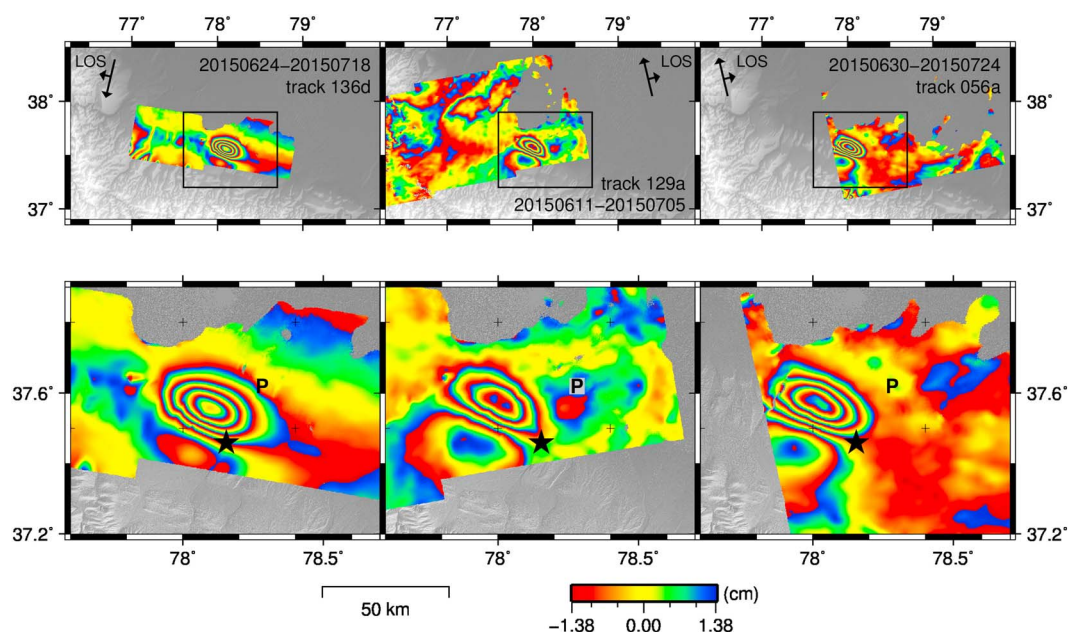


Figure 3. Coseismic interferograms from three Sentinel-1 tracks. (top row) The full extent of each interferogram used in this study, overlaid on a hillshaded topographic map. The lobe with the greatest change shows decreased range, that is, displacement toward the satellite. The time span of each interferogram is printed on the right, format *yyyymmdd*. Black box shows the area enlarged in Figure 3 (bottom row). Black stars show epicenter (USGS), P marks Pishan town.

3.2. InSAR

3.2.1. Coseismic Sentinel-1 Interferograms

We used InSAR data showing ground displacements to add an independent estimate of the earthquake centroid parameters and made use of the interferograms' dense near-field spatial coverage to constrain the fault location, slip extent, and slip distribution. The coseismic ground deformation was measured by constructing interferograms using data from the European Space Agency's (ESA) Sentinel-1A satellite, which carries a C-band synthetic aperture radar (SAR). We made interferograms from three tracks: two ascending and one descending (Figure 3 and Table 2), each had a time span of 24 days.

We processed the single look complex (SLC) images using the GAMMA software, following the method detailed in Appendix A. The final interferograms were geocoded and downsampled to 100 m. They show a simple two-lobed signal that varies smoothly over ~ 45 km, consistent with a blind event associated with buried faulting (Figure 3). The northern lobe indicates motion toward the satellite, with a peak line-of-sight (LOS) displacement of 11.8 cm for track 056a. The southern lobe represents motion away from the satellite,

Table 2

Details of the Sentinel-1A Interferograms Used in This Study

Track	Direction (asc/dsc)	Date 1 (yyyy/mm/dd)	Date 2 (yyyy/mm/dd)	$\perp B^a$ (m)	ΔT^b (days)	ΔPT^c (days)	Gradient removed ^d (cm/km)
<i>Coseismic Interferograms</i>							
129	asc	2015/06/11	2015/07/05	−39	24	2	−
056	asc	2015/06/30	2015/07/24	53	24	21	−
136	dsc	2015/06/24	2015/07/18	−26	24	15	−
<i>Postseismic Interferograms</i>							
129	asc	2015/07/05	2016/01/13	−46	192	−	2.43
056	asc	2015/07/24	2016/02/01	−133	192	−	4.19
136	dsc	2015/07/18	2016/01/02	19	168	−	2.72

^aPerpendicular baseline at scene center. ^bTime span. ^cPostseismic interval. ^dTopographically correlated phase gradient removed from the interferogram.

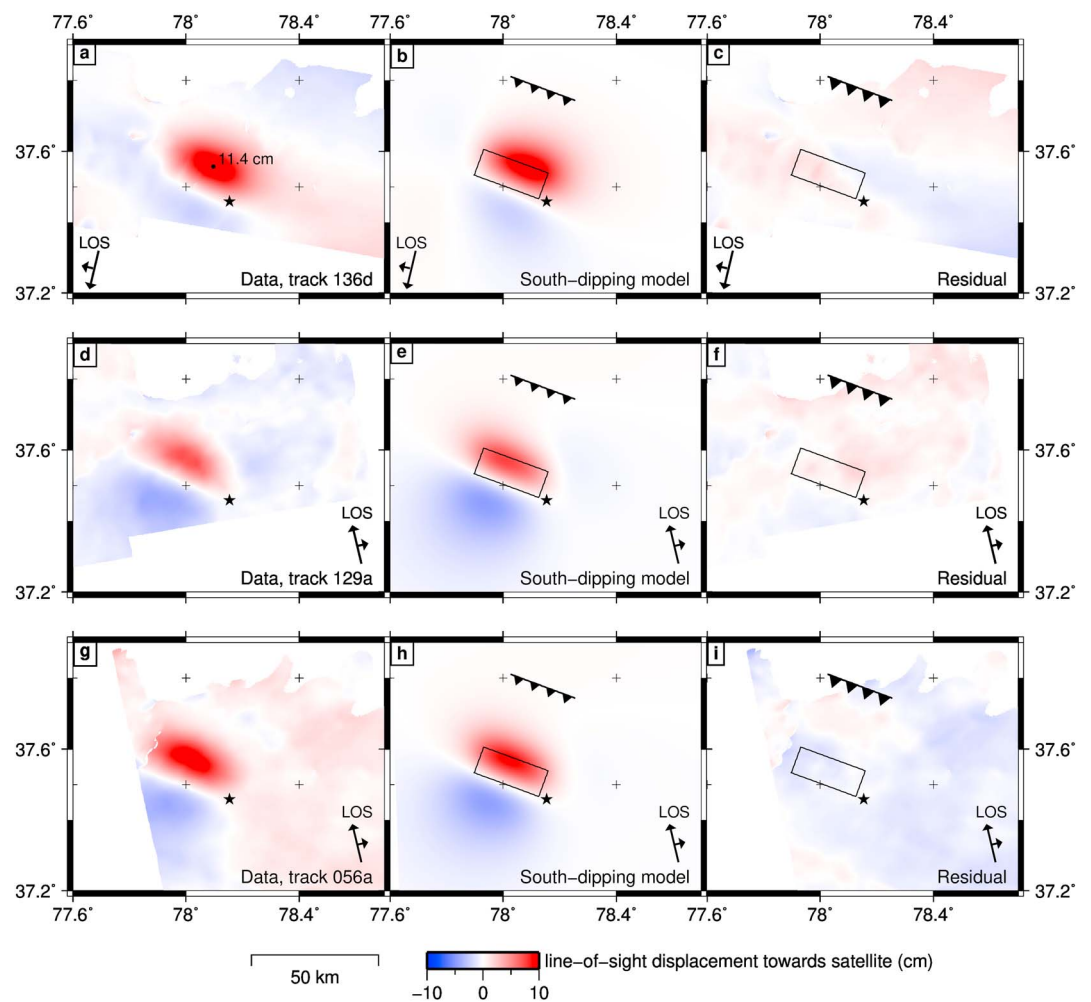


Figure 4. Results of uniform-slip models for the coseismic interferograms. (a) Unwrapped interferogram for track 136, location of maximum line-of-sight displacement toward the satellite of 11.4 cm is marked by the black dot. (b, c) Model and residual line-of-sight displacement for a south dipping fault model. Black ticked lines: updip projection of the fault plane to the surface. Black boxes: location of the uniform slip patch at depth. Black stars: epicenter (USGS). The equivalent figures for tracks (d–f) 129 and (g–i) 056.

with peak LOS motion of 5.9 cm for the same track. The three tracks cover the spatial extent of the entire deformation zone, although coherence is lost in some areas of the desert 50 km to the north of Pishan. While the two ascending interferograms covering the earthquake have the same look geometry, the different locations of the peak displacement of the earthquake relative to the SAR scene extents provides differing sensitivities to the ground deformation: the earthquake lies in the near-range portion of track 056a, with incidence angles $\sim 30^\circ$, indicating a greater sensitivity to vertical motion than east-west; for track 129a the incidence angle is nearer 45° in the far range, resulting in an equal sensitivity to vertical and east-west ground displacement.

3.2.2. Uniform-Slip Model

By fitting the ground displacements measured by InSAR with an elastic dislocation model (Okada, 1985; Wright et al., 1999) we obtained an independent estimate of the fault parameters. Our model comprises uniformly distributed slip on a rectangular dislocation in an elastic half-space, and we inverted for strike, dip, rake, location in x and y , top depth, bottom depth, length, and slip using the method described in Appendix B. The result is shown in Figure 4 and Tables 1 and 3. The focal mechanism agrees with the seismological results in section 3.1—it has a strike of 112° , 83° rake, and dip of 21° . Based on this mechanism, we believe the south dipping nodal plane is the fault plane because it agrees with the seismic reflection profile presented in Li et al. (2016) and section 4, whereas the north dipping plane would cross unbroken reflectors on the reflection profile and would give a poor fit to the postseismic displacements in section 6.

Table 3
Misfits Between our Models and Interferograms

Period	Dip direction	Uniform/Distributed	Track RMS (cm)			Overall RMS (cm)
			129	056	136	
Coseismic	South	Distributed	0.58	0.61	0.52	0.57
	South	Uniform	0.59	0.66	0.53	0.60
Postseismic	South	Distributed	0.38	0.30	0.37	0.36

3.2.3. Distribution of Coseismic Slip

To enable a more detailed comparison between coseismic slip and the long-term geological structure visible in seismic reflection profiles (section 4), we invert for the distribution of slip on the fault plane based upon the geometry found in section 3.2.2.

Taking the location, strike, dip, and rake from the uniform slip inversion, we extended the fault plane laterally, to 20 km depth and to the surface. The fault was split into patches of size 2 km in length and ~ 2 km in downdip width, and we carried out a nonnegative least squares inversion to solve for slip on each patch and for the interferogram nuisance parameters (Funning et al., 2005; Wright et al., 2004). In order to find a physically realistic solution, the method imposes a positivity constraint and Laplacian smoothing. The multiplying factor of the Laplacian (i.e., the strength of the smoothing) determines the balance between roughness and root-mean-square (RMS) misfit to the downsampled InSAR data. We chose a smoothing factor in the apex of the RMS misfit versus roughness curve, aiming to give the highest resolution possible without introducing spurious slip patches (Figure S1 in the supporting information). By inspection of the results from a variety of smoothing factors, we see that our main conclusions are not strongly influenced by the choice of smoothing factor from within a reasonable range.

Our preferred slip distribution is shown in Figure 5. It shows a single slip patch with a peak slip of 1.1 m, length of ~ 20 km, and downdip width of ~ 10 km. This model corresponds to the majority of the slip being at 9.5–13 km depth. We found no substantial slip above 8 km depth. Our interferograms have good coverage above the surface projection of the fault planes, and we can be confident in our conclusion that the earthquake was blind. Our results are consistent with field observations that found no evidence of surface ruptures (Lu et al., 2016) and are also consistent with our interpretation of local geomorphology (section 5), with seismic reflection profiles (Li et al., 2016) and the slip distributions of Zhang et al. (2016) and Sun et al. (2016) but contradict earlier InSAR-derived slip distributions by Wen et al. (2016) and He et al. (2016). The latter slip distributions each show ~ 20 cm of slip in the shallow portion of the fault plane, including at the surface.

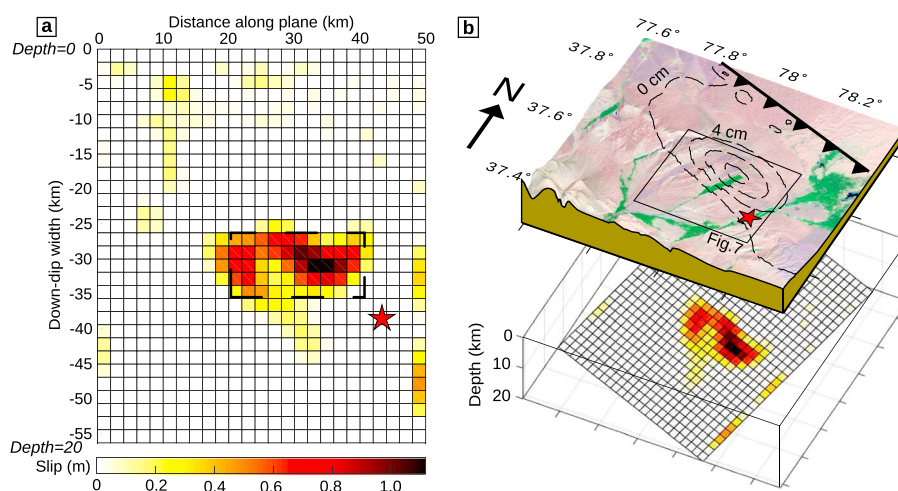


Figure 5. (a) Coseismic slip distribution. Black dashed lines show the extent of the fault plane from uniform slip model. Red star: hypocenter (USGS), projected onto the fault plane. (b) Three-dimensional perspective view of the slip distributions and the fault's surface projection. Black dashed lines: contours of the models' vertical ground displacements plotted in 4 cm increments from zero. Black box: location of Figure 7. Red star: epicenter (USGS).

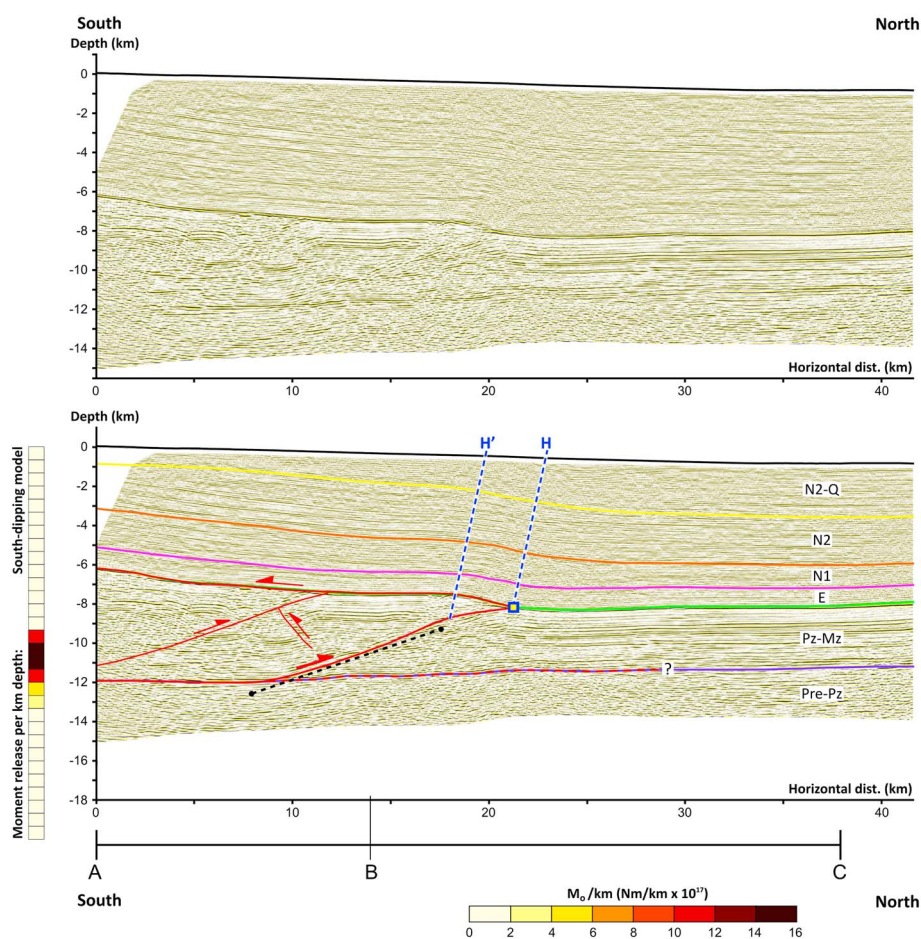


Figure 6. Seismic reflection profile along the line A-B-C (Figure 1) (adapted from Li et al. (2016) and with our modified velocity model). (top) Uninterpreted. (bottom) With interpretation. The back thrust at 10 km horizontal distance is our interpretation, all other interpretation is taken from Li et al. (2016). Dashed black lines: fault planes found by our uniform slip models. H' and H: fold axes. N2-Q: Pliocene-Quaternary 3.6 Ma to present (Zheng et al., 2000). N2: Pliocene. N1: Miocene. E: Paleogene. Pz-Mz: Paleozoic-Mesozoic. Pre-Pz: Precambrian. (left) Colored squares show the depth distribution of moment release per unit depth in our finite slip model.

Our interferograms show that there is no discontinuity at the surface projection of the fault, and we suggest that the shallow slip suggested by other studies (He et al., 2016; Wen et al., 2016) is an artifact as a consequence of their incomplete data coverage in the Taklamakan Desert (Figure 1) above the surface projection of the fault plane—the area most sensitive to shallow slip.

4. Comparison of Fault Positions Derived From InSAR and Seismic Reflection

A high-resolution seismic reflection profile that crosses the coseismic slip patch has been published and described by Li et al. (2016). Their profile D-D' is a combination of two individual profiles that overlap. Here in Figure 6 we show a modified version (see below) of the more northerly of the two, which crosses the whole of the coseismic fault and was taken along line A-B-C (Figure 1). The Li et al. (2016) profile shows a stacked series of northward vergent thrusts forming a wedge between a basal detachment at the Paleozoic-pre-Paleozoic boundary and a roofing detachment at the Cenozoic-Mesozoic boundary. Above this wedge lie ~7.5 km of Cenozoic sediments which are folded into a monocline above the wedge tip and are not offset by any significant faults (Figure 6). The wedge is thickened by up to a factor of 2 compared to the undeformed Paleo-Mesozoic unit north of the thrust tip. This thickening, the subhorizontal dip of reflectors within most parts of the wedge and the smooth, simple folding of the roofing detachment suggests that, instead of forming by rotation of coherent blocks, internal deformation of the wedge has played a part in its development in addition to slip on the bounding faults. We have added a small back thrust to the interpretation

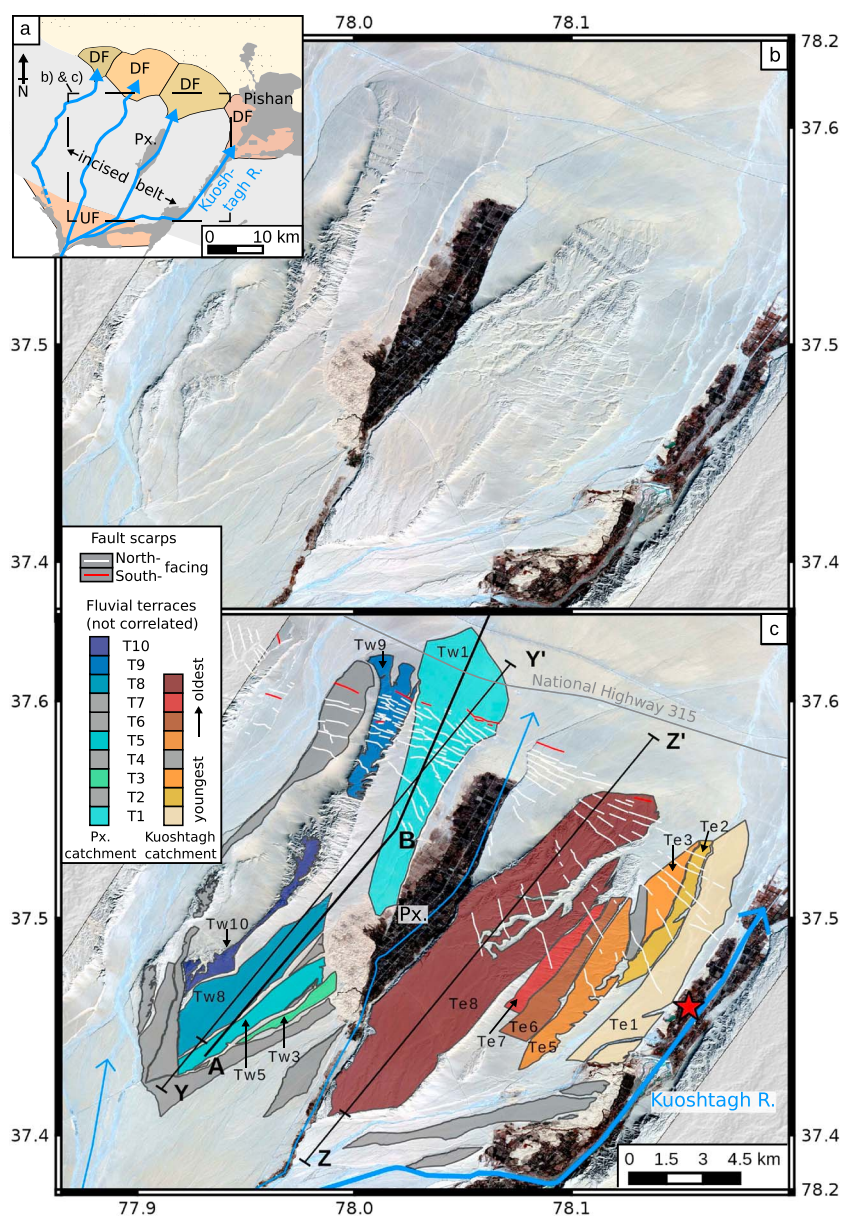


Figure 7. (a) Simplified sketch map of the geomorphology around the Pishan terraces. UF: fan upstream of the fold. DF: downstream fans. (b) Uninterpreted SPOT7 imagery of the Pishan terraces and surface scarps. A 1.5 m resolution panchromatic imagery has been used to pansharpen 6 m resolution multispectral (RGB) imagery. (c) Interpreted version of the same area. Colored polygons: terraces with profiles in Figure 8. Solid black lines: profiles Y-Y' and Z-Z'; distance markers at 0 km. Red star: epicenter (USGS). Px: Pinxinaxiang town. Thick blue arrow: active river. Thin blue arrows: inactive drainage routes. Line A-B: part of the seismic reflection line shown in Figure 1. Grayed-out terraces were not used in Figure 8.

by Li et al. (2016). The back thrust is the most prominent sign of internal deformation within the wedge and offsets a prominent horizon at 10 km horizontal distance. It is not enough alone to account for the wedge thickening, but it could contribute to keeping the reflectors above it subhorizontal; if the back thrust slipped 1 m for every 2.2 m of slip on the frontal ramp, then the two faults would have equal throw. The back thrust is in the hanging wall of the ramp and presumably is transported by it. Due to the constant thickness of the Pliocene unit (N2 on Figure 6), Li et al. (2016) label it as pregrowth stratum and conclude that the initiation age of the deformation of the wedge is ≤ 3.6 Ma.

We note a very good agreement in location, but a difference in dip and depth between our InSAR-derived fault plane estimate and the frontal wedge thrust in Li et al.'s (2016) reflection profile interpretation. However,

Li et al. (2016) used a first-approximation velocity model: a uniform velocity of 3,000 m/s. While 3,000 m/s is appropriate for the Cenozoic section, it is likely to be an underestimate for the deeper units. The Paleozoic and Mesozoic units are predominantly sediments, such as limestones, dolomite, and gypsaceous mudstone, whereas the pre-Paleozoic units consist of metamorphic and volcanic rocks (Lu et al., 2016). Data from Kashgar, where the lithologies are thought to be similar to the Pishan area, indicate that the velocity of the pre-Cenozoic units is commonly in the range $\sim 4,500$ – $5,500$ m/s, whereas the Cenozoic unit is usually $\sim 2,800$ – $3,200$ m/s (Heermance et al., 2008). We have therefore updated the seismic reflection profile, changing the time-depth conversion to incorporate a velocity model of 3,000 m/s for the Cenozoic sediments and 5,000 m/s for earlier units. The result is shown in Figure 6, and the location of the profile is shown in Figure 7. The top four units (N2-Q to E) of the newly converted profile are unchanged from the original, but below that, the effect is that reflections have been deepened, with the greatest effect on the deepest parts of the profile. The Precambrian-Paleozoic unit boundary level with the zero kilometer horizontal distance marker on Figure 6 has increased in depth from 10.5 km to 12 km and the dip of the front ramp fault has increased from 10 – 14° (Li et al., 2016) to 16 – 20° , bringing the dip into line with teleseismic and InSAR estimates for the 2015 earthquake (Table 1).

The southward dipping InSAR solution matches the independently constructed updated seismic reflection profile closely in depth, location, and dip. We note that, as always, both the InSAR fault location and geometry and the reflection profile have uncertainties associated with them and it would be unwise to rely on interpreting the fine details of the comparison. However, the excellent match is strong evidence that the causative fault of the Pishan earthquake was the frontal wedge thrust that is visible in the reflection profile (Li et al., 2016). This consistency indicates that InSAR solutions of active faulting could provide useful constraints on the velocity migration of seismic reflection profiles, particularly in areas of thick sedimentary cover lacking borehole constraints on seismic velocities.

5. Folding in the Geomorphology

5.1. Fluvial Terrace Deformation

The Pishan earthquake occurred under the alluvial apron that skirts the southern edge of the Taklamakan Desert. Drainage runs to the north-northeast from the Western Kunlun Mountains. A belt of incision ~ 25 km wide crosses from northwest to southeast parallel to the strike of the fault (Figure 1, section 2). In some places it exposes pre-Quaternary sediments (Li et al., 2016; Matte et al., 1996), and in the Pishan area it preserves a sequence of folded fluvial terraces. The mainshock epicenter and the surface-projected fault plane lie in this terraced area (Figures 1 and 5b).

We mapped the terraces adjacent to the Kuoshtagh and Pinxinaxiang town rivers using 1.5 m resolution SPOT7 optical imagery, the 30 m SRTM1 DEM, GoogleEarth, and Bing Maps (Figure 7). We inferred relative age from mapping, surface preservation, and height and labeled the terraces in order of increasing age but were not able to correlate terraces between the two catchments (i.e., Te1 does not necessarily correlate with Tw1). Topographic profiles projected onto lines Y-Y' and Z-Z' are shown in Figure 8. In order to isolate the tectonic contribution to topography, it is necessary to remove the depositional slope. We have assumed and removed a linear depositional gradient of 0.92° as described in the next paragraph. After the removal of this depositional gradient, we assume that any remaining elevation changes on individual terraces are due to postdepositional tectonic deformation (e.g., Molnar et al., 1994; Poisson & Avouac, 2004; Walker et al., 2015).

Quantification of the original depositional gradient is not straightforward. Upstream, a single river forms one broad fan as it leaves the mountains (UF on Figure 7a), but through the incised belt drainage has been split into four channels, each of which has formed its own small fan downstream of the terraces (DF on Figure 7a). This prevents us from following coeval depositional surfaces upstream and downstream of the deformed terraces with confidence. We have therefore taken the depositional gradient from only the upstream fan and have removed it from all our profiles in Figures 8c and 8d. We note that alternative options, such as using the gradients of the channels or of the downstream fans, lead to different assumed depositional gradients in the range ~ 0.70 – 0.95° , but each of these gradients are potentially contaminated by later sediment deposition or by slopes perpendicular to the profile lines. Our conclusions would be unchanged by using a depositional gradient from anywhere within the ~ 0.70 – 0.95° range.

The profiles in Figures 8c and 8d show an anticline with a width of 15–20 km. Higher, older terraces show greater magnitudes of folding than the lower, more recently abandoned ones, showing that the fold growth

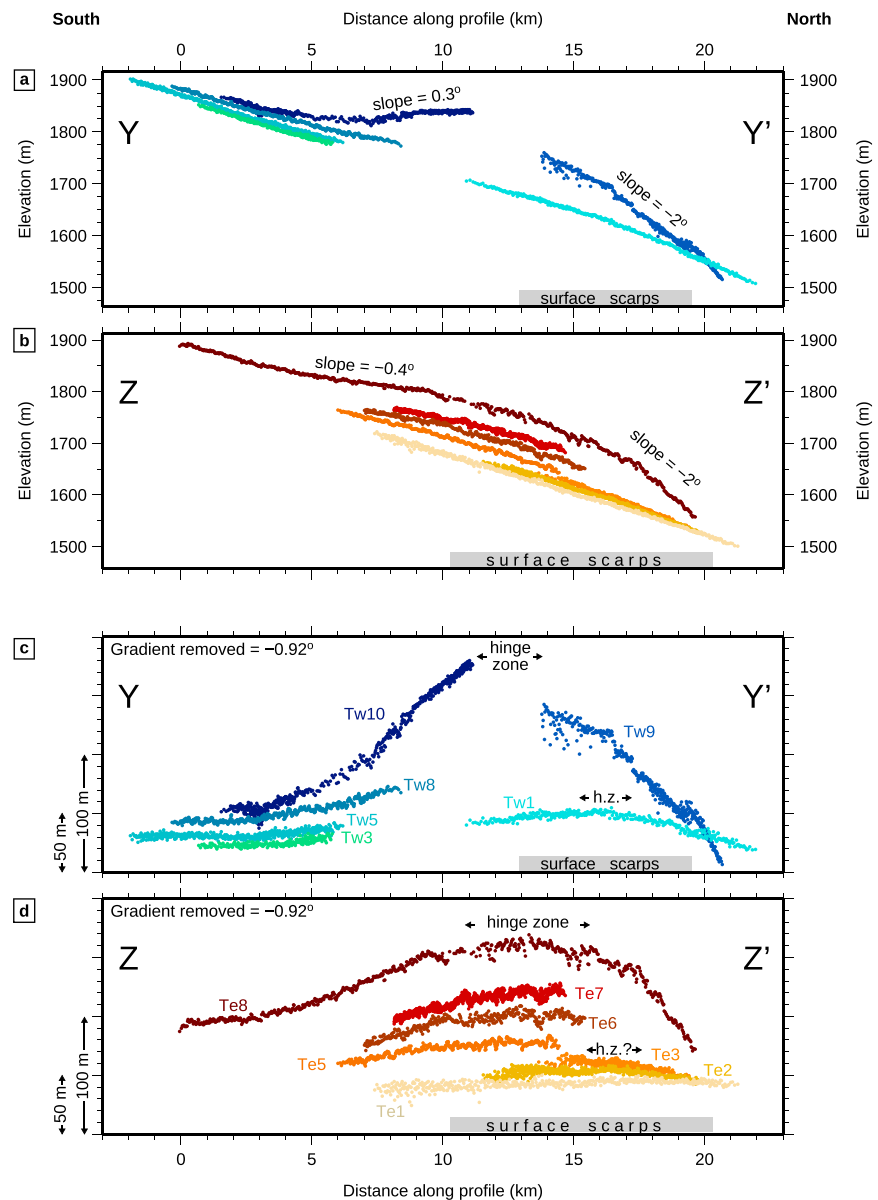


Figure 8. SRTM1 profiles of the colored terraces in Figure 7. (a) Profiles from the western terraces showing true elevations and vertically exaggerated true slopes. (b) Profiles from the eastern terraces, true elevations, vertically exaggerated true slopes. (c and d) The profiles from Figures 8a and 8b after the removal of the gradient of the fan upstream.

has been ongoing over the period during which the terraces were formed. The main fold crest is at 11–14 km horizontal distance on our profiles, but the lowest terraces (Tw1 and Te1–3) have a slightly more northerly crest at around 16 km, suggesting that basinward hinge migration may have taken place.

Our remote mapping found no single primary surface scarp or rupture and no surface evidence for a major fault reaching the surface. This is consistent with the seismic reflection profile (section 4, Li et al., 2016) which shows continuous, unbroken, reflectors above and updip of the coseismic fault. There is, however, a series of surface scarps distributed over 5–10 km across the steep northern limb of the fold (Figures 7 and 8). The scarps cut across fluvial terrace risers, strike roughly parallel to the fold axis, and are larger on older terraces. This shows that the scarps are tectonic in origin and, like the fold, their growth has been ongoing for some time. We find local instances where the scarps are nonparallel, bend, and truncate each other. This pattern is not compatible with being bedding-plane slip. Unlike conventional bending moment faults, the scarps are asymmetrical. The great majority are south up, although we do find a small number which are north up (Figure 7).

The scarps are not spread across the fold crest, and there are none on the southern limb. By summing the scarp height measurements made by Li et al. (2016) (T3 on their Figure 2b) we estimate that one half to one third of the height change on Tw9 is provided by the scarps, which are each 1–20 m high.

5.2. Abandonment Age of the Terraces in the Hotan Fold and Thrust Belt

In section 5.1 we showed that uplift has been ongoing with a consistent pattern since at least the formation of the oldest terraces in Figure 7, and it is useful for the following discussion to place bounds on that duration. The ages of the terraces at Pishan have not been directly determined, but as the formation of terraces is likely to be climatically controlled, we can estimate a likely age range from terrace abandonment dates from elsewhere, provided that the two places have a similar climatic history over the relevant time period. Dating from aggradational terraces around the Tarim Basin and the Qilian Shan, 250–800 km distance from Pishan, has found ages from a few thousand years to >140 ka (e.g., Avouac & Peltzer, 1993; Li et al., 2013, 2015; Mériaux et al., 2004; Saint-Carlier et al., 2016) and slightly farther afield Stockmeyer et al. (2017) recently dated a suite of terraces on the northern front of the Tien Shan as 5–212 ka. Given that present-day GPS velocities could allow no more than 3 mm/yr in the strike-perpendicular direction (Ge et al., 2015) and that the oldest of the terraces (Tw10 and Tw9) are deformed vertically by at least 120 m (Figure 8), it is likely that the oldest Pishan terraces are toward the older end of these age ranges.

6. Postseismic Deformation

6.1. Postseismic Interferograms From Sentinel-1

We constructed three postseismic interferograms covering the time from 2 days to 7 months after the earthquake (Table 2) using the same three tracks as used in section 3.2, including both ascending and descending viewing geometries. We processed the interferograms using the methods described in section 3.2.1 and found that coherence is similar to that of the coseismic interferograms.

Topographically correlated signals in the interferograms partially obscure the tectonic signal. This made it necessary to add the additional processing step of atmospheric correction for the postseismic interferograms. After unwrapping we made an empirical linear correction, which is a widely used and tested method for atmospheric correction (e.g., Bekaert, Walters, et al., 2015; Copley et al., 2015; Elliott et al., 2008). After unwrapping we worked only on a subset region of interest from the large Sentinel-1A TOPS mode interferograms. We used cropped areas (displayed in Figures 9a–9c) that are nonetheless larger than a scene from traditional earlier SAR instruments, to avoid any chance of incorporating areas with substantially different atmospheric conditions into our single empirical fit of elevation versus phase (a topic discussed by Barnhart and Lohman (2013) and Bekaert, Hooper, et al. (2015)). There are no topographic barriers or bodies of water in the region that would make it particularly susceptible to spatially variable atmospheric conditions, and a spatially invariant linear correction is sufficient for our purposes. The elevation-phase correlation coefficients for tracks 056, 136, and 129 were 0.9, 0.8, and 0.4, respectively (Figure S2). We then proceeded to the flattening step. There is some coregistration error in t129a which causes apparent motion away from the satellite at burst boundaries (regularly spaced lines parallel to LOS on Figure 9b) but our conclusions are the same if t129a is excluded.

Figure 9 shows the postseismic displacement signal. Like the coseismic signal it consists of two lobes with the same sign of displacement as for the earthquake deformation but about one fifth of the magnitude. The northern, more prominent lobe indicates motion toward the satellite, with a peak LOS displacement of 2.7 cm in track 056a. A small-displacement lobe of motion away from the satellite is present farther south but is less well resolved above the noise. The signal is asymmetrical for all three tracks: profiles through the unwrapped interferograms show it has a steep northern (updip) side and a broader southern limb that is punctuated by a change in gradient (at ~12 km in Figures 9d–9f), south of which the displacement gradient is decreased. The axis of peak postseismic displacement is shifted approximately 1 km NNE (in the updip direction) from the coseismic peak axis. The peak motion away from the satellite is shifted by ~7 km in the downdip direction from the coseismic peak and the transition from uplift to subsidence is also shifted southward in the postseismic interferograms.

6.2. Causative Mechanisms of the Postseismic Displacements

Postseismic surface displacement can be caused by a number of mechanisms. In this case, the earthquake's moderate magnitude, midcrustal slip depth, and the short timescale and length scale of the observed signals rule out mantle relaxation as a cause of the signal, which is typically observed for much larger events (Wright et al., 2013). The polarity of the postseismic and coseismic displacement patterns is the same. This is not

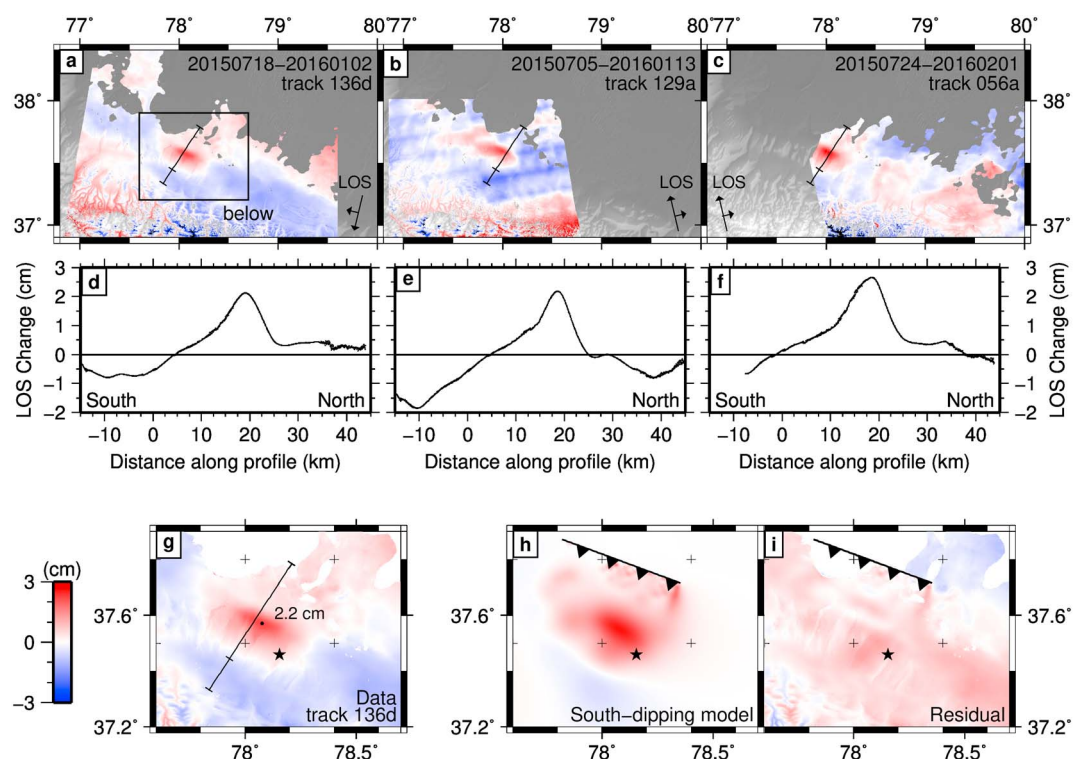


Figure 9. (a–c) Unwrapped postseismic interferograms as detailed in Table 2. Timespan is shown in the top right of each panel, format yyyy/mm/dd. (d–f) Profile along the line marked on Figures 9a–9c. (g–i) Data, distributed slip models, and residuals for the postseismic deformation signal for track 136d. Black, ticked lines show surface projection of the fault planes shown in Figure 10, with ticks pointing downdip. Black star shows epicenter (USGS).

consistent with poroelastic rebound being the dominant postseismic deformation mechanism (Jónsson et al., 2003), although we cannot discount it playing a minor role. In situations such as this, where an unfaulted pile of sediments overlies a blind thrust fault (section 4), Lyzenga et al. (2000) identified anelastic relaxation of the sediments as another potential mechanism for postseismic relaxation. This would not explain the change in gradient in the southern limb in the profiles (section 6.1 and Figure 9) but could be acting in combination with another mechanism. Afterslip is the final candidate for the cause of the postseismic displacements. We cannot rule out a combination of mechanisms and, given the asymmetrical and kinked profile discussed in section 6.1, a combination may be likely.

6.3. An Afterslip Model for the Postseismic Interferograms

Using the fault plane extent, rake, slip patch size, and smoothing factor taken from our models of coseismic slip distribution (section 3.2.3), we applied the same method to model the distribution of afterslip, assuming for now that afterslip is the sole tectonic cause of the postseismic displacements in Figure 9. The resulting slip distribution has the shape of a ring that sits around the base and sides of the coseismic slip patch and partially overlaps with it at the top (Figure 10). The ring shape arises from the short wavelength peak in the displacements superimposed on a longer wavelength, which was discussed in section 6.1 and Figure 9. Slip at two distinct depths is required to match the kink (section 6.1) and is seen across a range of different model smoothing factors (Figure S3). This distribution is resolvable above the noise level: Figure 10b shows the standard deviation from 100 interferograms that were perturbed (as in section 3.2.2) and then modeled using our preferred smoothing value. It is also not an artifact of the inversion: in tests where the LOS ground displacements produced by synthetic filled or ring-shaped slip patches were downsampled to the same resolution as used in the slip inversion and then modeled, the ring and filled shapes were recovered (Figures S4 and S5).

Afterslip distributed around and slightly overlapping the coseismic slip patch is in line with theoretical models including but not limited to the case where a velocity-weakening patch is surrounded by velocity-strengthening fault plane (Avouac, 2015; Helmstetter & Shaw, 2009; Marone et al., 1991) and has already been observed in some cases (e.g., Liu et al., 2016; Miyazaki et al., 2004; Thomas et al., 2014). In this case our

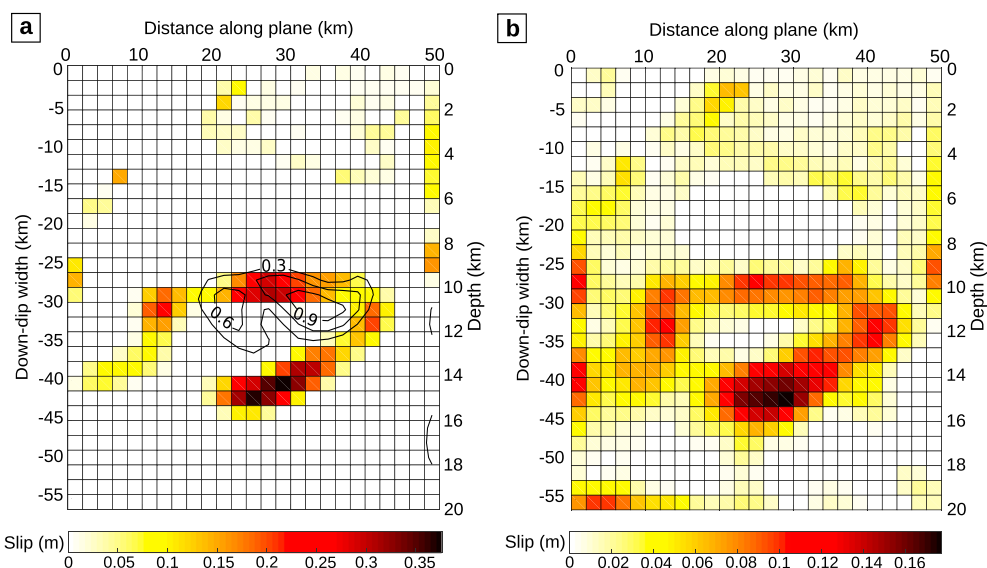


Figure 10. (a) Slip distribution for an afterslip model on a planar, south dipping fault showing contours of coseismic slip in units of meters. (b) Standard deviation of the slip distribution at our preferred smoothing value.

interpretation is limited by the similarity of surface displacement patterns produced by different models. Lyzenga et al. (2000) showed that the surface displacement due to anelastic deformation of overlying sediments can be well fit by a model of slip on a discontinuity close to the top half of the coseismic slip plane. Although the plane is predicted to be slightly above and steeper than the coseismic plane, these differences are too small and our ability to isolate the shallow part of the signal is insufficient for us to confidently attribute the signal to one cause or the other. Similarly, we were not able to distinguish between a planar fault geometry and a fault that flattens to a dip of 1° below 12 or 12.7 km depth at the base of the coseismic slip patch. They fit the data equally well, and although the nonplanar slip models are more disjointed than the planar model, they are more consistent with the thrust system interpreted from the seismic reflection profile discussed in section 4, which shows the fault shallowing into a flat at a depth of around 12 km (depending on the seismic velocity assumptions). The 12 geometries tested all showed afterslip at at least two distinct depths, as in Figure 10.

7. Pishan Earthquake in the Context of the Tarim-Tibet Margin

Low-angle thrust earthquakes of the type seen at Pishan are relatively rare on the continents, with most shortening happening on faults dipping at 30° – 60° (Sibson & Xie, 1998). In order to investigate the possible causes of this unusual dip angle, we need to investigate the material properties of the lithospheres of the Tarim Basin and Tibetan Plateau.

An M_w 5.7 earthquake struck the central Tarim Basin on 8 March 2012 (Figure 1, inset). We have modeled this earthquake using the same seismological method as described for the Pishan earthquake (section 3.1), and the results of our waveform modeling are shown in Figure 11. This thrust-faulting earthquake occurred on an E–W to SE–NW oriented plane, slightly to the north of the Bachu uplift (Figure 1, inset), a thrust-bounded basement uplift in the central Tarim Basin (e.g., Allen et al., 1999; Tong et al., 2012). The depth of the event is well constrained at 51 ± 5 km due to the clear separation in time between the direct wave and the depth phases (Figure 11; calculated using the Tarim Basin velocity model of Huang et al. (2017)), possibly suggesting that the lower crust of the region is anhydrous (e.g., Jackson et al., 2008). The earthquake depth is within error of the depth of the Moho (Zhang et al., 2011), and it is not possible to determine whether it was in the lower crust or upper mantle. Published depths for other earthquakes in this region have shown that brittle deformation spans the full thickness of the crust down to at least 44 km (Huang et al., 2017).

The forces transmitted between mountain ranges and their associated lowlands can be calculated by summing the lateral differences in vertical normal stresses between the two lithospheric columns (e.g., Artyushkov, 1973; Dalmayrac and Molnar, 1981). We have done this for Tarim and northern Tibet using

Tarim Basin, 8/3/2012

Minimum Misfit Solution:
 Depth: 51.0 km
 Magnitude: $M_w = 5.7$
 Moment: 4.321×10^{17} N m
 Mechanism: Strike: 125°
 Dip: 41°
 Rake: 116°

Source-Time Function
 0 50
 Waveform Timescale

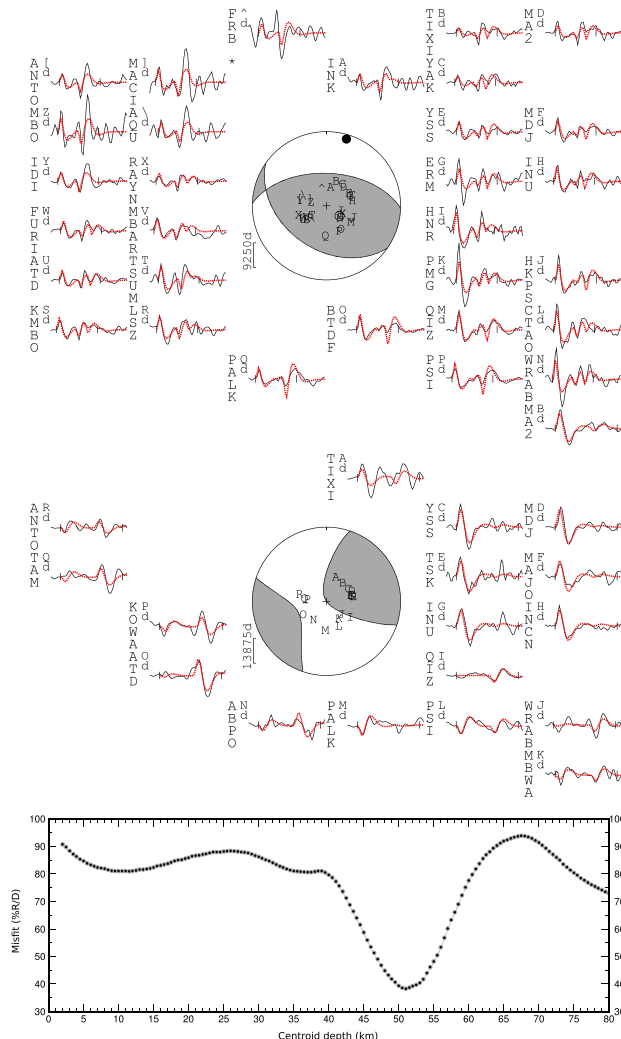


Figure 11. As in Figure 2 but for the 8 March 2012 M_w 5.7 Tarim Basin earthquake.

the method of Copley and Woodcock (2016), which includes the effects of crustal and lithospheric thickness contrasts, and thermal and chemical effects on the density. Using this method, we estimate that the Tarim Basin and northern Tibetan Plateau exert upon each other a force of $4 \pm 2 \times 10^{12}$ N per meter along strike (with the error estimate encompassing our uncertainties in the thermal structure of the mountains and lowlands). We can resolve this total force onto the faults that have broken in earthquakes in the Tarim Basin and, by doing so, estimate an upper limit on the vertically averaged shear stress on the faults of 40 ± 20 MPa. If the faults were able to support more shear stress than this value, there would be insufficient force available to make them break in the earthquakes we have observed. This estimate is an upper limit, because it assumes that all stresses are supported in the brittle, rather than ductile, part of the lithosphere. Our estimated shear stresses correspond to an upper limit on the effective coefficient of friction of 0.05 ± 0.025 . Although much lower than laboratory-derived estimates of fault friction (e.g., Byerlee, 1978), our estimates are within the range calculated for other deforming zones (e.g., Copley et al., 2011; Herman et al., 2010; Lamb, 2006) and require high pore fluid pressures, intrinsically weak material on fault planes, or both.

We now compare the Tarim Basin with the Tibetan Plateau. Within the plateau, earthquakes are confined to the upper 15 km of the crust, along with isolated deep events close to the Moho related to the underthrusting of Tarim from the north and India from the south (Craig et al., 2012). Unlike in the Tarim Basin, the majority of the crust deforms by ductile flow. We can balance the forces exerted across the margin (Figure 12) to estimate the viscosity of the Tibetan lithosphere. A crucial input into this procedure is the stress drop in the Pishan earthquake. The full procedure is described in Appendix C. There is a trade-off between the values of η_u (the average viscosity of the upper part of the ductile layer) and η_l (the average viscosity of the lower part of the ductile layer) on Figure 12. If they are taken to be equal, then $\eta_u = \eta_l = 9 \pm 5 \times 10^{21}$ Pa s. If they are assumed to be different by 1 or 2 orders of magnitude, then either $\eta_u = 1.2 \pm 0.6 \times 10^{22}$ Pa s and $\eta_l = 1.2 \pm 0.6 \times 10^{21}$ Pa s or $\eta_u = 1.2 \pm 0.7 \times 10^{22}$ Pa s and $\eta_l = 1.2 \pm 0.7 \times 10^{20}$ Pa s.

The analysis in this section has highlighted a stark contrast between the rheology of the Tarim Basin and the Tibetan Plateau. Tarim deforms by brittle failure throughout its thick seismogenic layer, whereas the majority of the Tibetan crust deforms by ductile creep with a viscosity of only 1 or 2 orders of magnitude higher than the convecting mantle. This contrast may represent

the cause of the unusual low dip of the Pishan earthquake fault plane. Areas with a large seismogenic thickness also have a large elastic thickness (Jackson et al., 2008) and will bend over long wavelengths. The upper surface of the Tarim Basin will therefore only dip gently as it underthrusts southern Tibet. The weak and ductile Tibetan lithosphere can flow over and around the underthrusting Tarim Basin crust, allowing the convergence between the two to continue. This strength contrast therefore leads to the presence of low-angle thrusting on the northern margin of Tibet. Similar strength contrasts are likely to be responsible for the low-angle thrusts on the southern margins of the Tien Shan and Tibetan Plateau (Allen et al., 1999; Elliott et al., 2016; Wang et al., 2011).

8. Discussion

8.1. Timing of Topographic Growth

The observations presented above allow us to place some constraints on the mechanisms controlling topographic growth at the northern Tibetan range front and its timing within the seismic cycle. Recent work on reverse faults has attributed topographic growth to various mechanisms depending on setting,

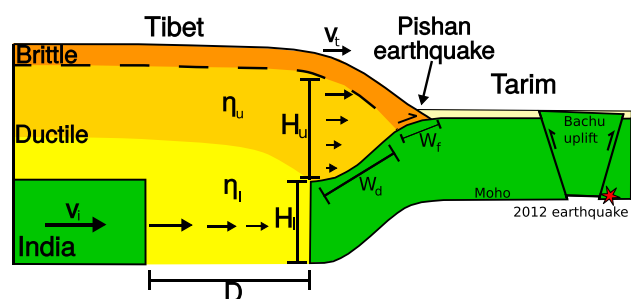


Figure 12. Schematic sketch of the deformation in northern Tibet and the Tarim Basin. V_i : rate of convergence between India and the Tarim Basin. V_t : rate of convergence between the western Tarim Basin and northern Tibet. η_u : average viscosity of the upper portion of the ductile Tibetan lithosphere. η_l : average viscosity of the deeper portion of the ductile Tibetan lithosphere. H_u , H_l , W_d , W_f , and D refer to the distances as shown.

including: coseismic slip (Le Béon et al., 2014; Stein & Ekström, 1992), post-seismic afterslip on a single fault (Copley, 2014; Elliott et al., 2015, 2016; Stein and Ekström, 1992; Zhou et al., 2016), afterslip on multiple faults (Copley & Reynolds, 2014; Mackenzie et al., 2016) and slip on faults or sections of fault that play an apparently minor part in the region's seismic activity (Mackenzie et al., 2016; Melnick, 2016; Whipple et al., 2016).

We have shown using InSAR and seismology (sections 4 and 3.1) that the Pishan earthquake took place on one of the main faults within the Hotan Fold and Thrust Belt and is the most basinward of a set of stacked thrusts. The lack of any major fault trace at the surface, the unbroken Cenozoic reflectors in the reflection profiles and our findings that slip in the Pishan earthquake stopped at depth, and there was no shallow (<9 km) afterslip, show that the Cenozoic sediments accommodate permanent strain in a distributed manner by folding. Our results allow us to discuss the deeper driving processes and show that deformation in the time between significant earthquakes is likely to contribute significantly to building

topography. Despite its location on a prominent fault, the coseismic ground displacements from the earthquake do not align with the anticline in Quaternary fluvial terraces (Figure 13). The peak coseismic uplift is offset approximately 5 km into the foreland from the topographic fold axis of the higher terraces and approximately 2 km from the axis of the lower, younger terraces. Although the northern part of the anticline was uplifted, the southern part subsided during the earthquake. The peak ground uplift in the first 7 months of postseismic motion is a farther kilometer offset from the topographic fold axis compared to the coseismic uplift (Figure 13).

Our results show that distributed postseismic relaxation of the Cenozoic sediments after the coseismic stress change also cannot explain the anticline growth. Lyzenga et al. (2000) showed that the displacement signal

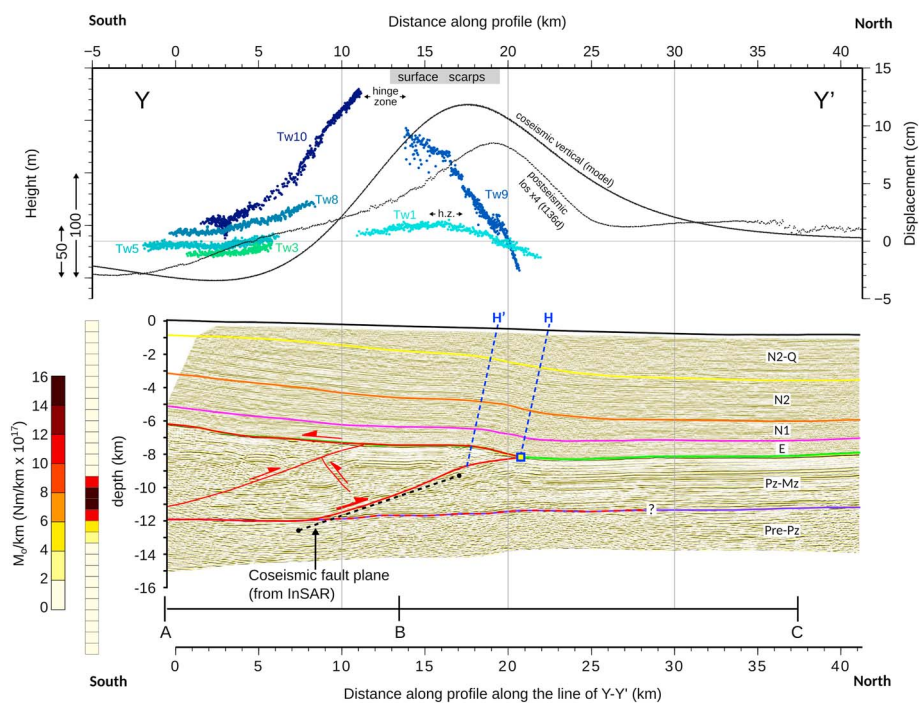


Figure 13. Aligned profiles along an extended Y-Y' line. Topographic profiles taken from Figure 8c are overlain on the vertical displacements from our coseismic model and a 4 times vertically exaggerated profile through the postseismic interferogram. (bottom) Seismic reflection profile duplicated from Figure 6. Black dashed line shows fault plane from uniform slip model. Colored sidebar with key shows coseismic moment release with depth in our distributed slip model. Note that no modification of the seismic reflection data was made to account for the 15° change in azimuth at B. There is no misalignment along A-B, where the profiles are parallel, but along B-C the seismic profile is stretched in comparison to the topography and InSAR. This results in ~230 m misalignment at 20 km horizontal distance and <1 km at 40 km.

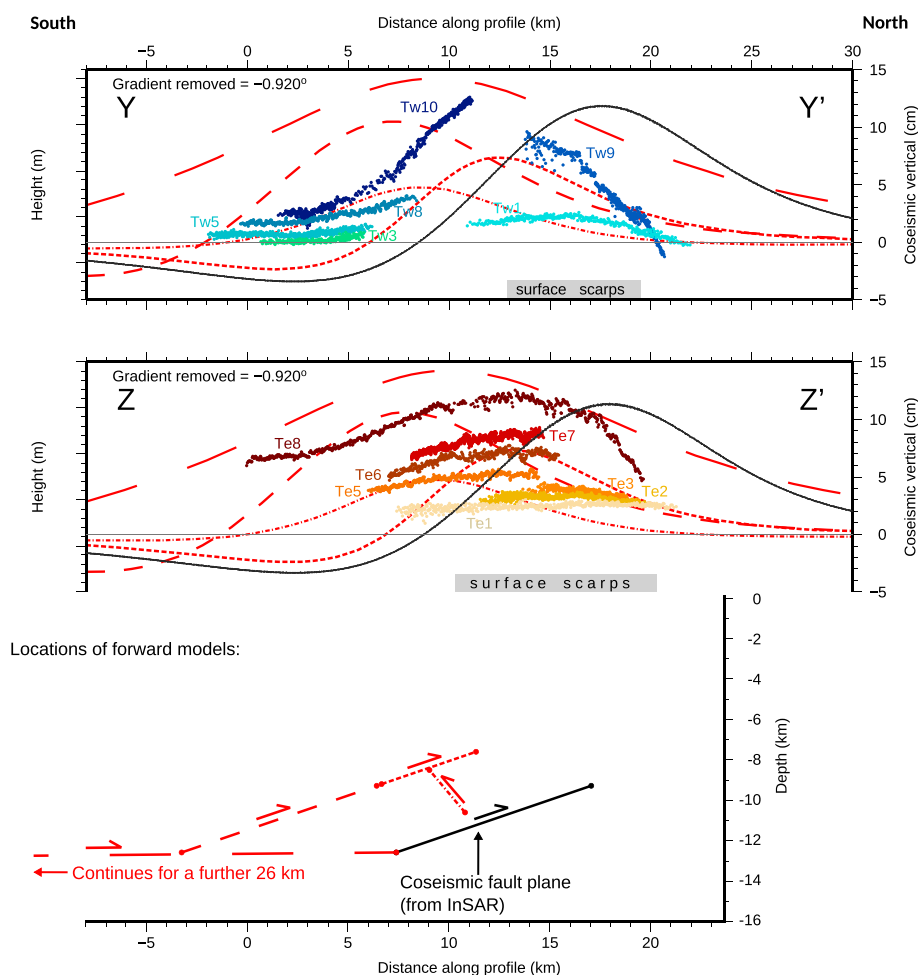


Figure 14. Profiles through coseismic uplift pattern and terraces as in Figure 13. Red dashed lines show the results of a suite of forward models, assuming rake 90, magnitude arbitrary. Long dashes: detachment, dip 0.5°, downdip length 46 km. Medium dashes: ramp to south of 2015 earthquake's fault, dip 21°, depth range same as coseismic. Short dashes: ramp to south, dip 21°, depth 7.6–9.2 km. Dot-and-dash: back thrust, dip 50°, depth 8.5–10.6 km.

caused by sediment relaxation after an earthquake would appear similar to afterslip close to the upper half of the coseismic fault plane, as we have in our afterslip-only postseismic models (Figure 10 and section 6). This would produce an uplift peak several kilometers to the north of the fold crest. We can therefore conclude that the anticline and minor scarps described in section 5 did not form coseismically by repetitions of the July 2015 earthquake.

The geometry shown in Figure 13 suggests the nature of the deformation in the time between 2015-like earthquakes. It shows a flat detachment leading to a ramp, which has little or no offset at its tip. In this configuration northward horizontal flux of material decreases from south to north, leading to material being forced upward into an anticline (e.g., Burbank & Anderson, 2012; Suppe & Medwedeff, 1990; Wickham, 1995). We have seen the effects of slip on the ramp in the 2015 earthquake. The remaining mechanisms are therefore slip on different parts of the ramp and flat (particularly interseismic transport on the detachment) and internal deformation of the wedge. Although only a few key faults have been interpreted in Figure 6, both the wedge geometry and the reflections show that it has undergone internal thickening in the past and may still be doing so.

The long wavelength part of the postseismic signal did uplift the anticline in the months following the earthquake. This signal is due to the afterslip below the coseismic slip patch that we discussed in section 6 and could be an important part of the ramp and flat slip. However, with <8 mm uplift of the fold crest during the earliest, fastest phase of postseismic motion, future studies will need to find a strongly sustained signal

for it to have been a major contributor to the ≥ 120 m anticline. Brief inspection of interferograms from early 2016 to early 2017 shows that the rate of postseismic displacement for that period was less than half of the rate for the earlier postseismic period shown in Figure 9 and that the peak signal continued to be north of the fold crest.

Figure 14 shows ground displacements from elastic forward models of earthquakes on the detachment, on the more loosely interpreted $\sim 21^\circ$ -dipping thrust fault 10 km behind the 2015 earthquake, and on the small back thrust. It is important to note that these models do not include the interseismic strain which loads the faults, and therefore, Figure 14 does not show the total ground deformation resulting from transport on these faults over the > 100 ka timescale recorded by the fluvial terraces. The uplift patterns differ from the terrace folding, and we are not able to account for the anticline by coseismic slip alone on any of these other mapped faults. This is not surprising because as we have stated, over the timescale of multiple earthquake cycles, interseismic and additional aseismic deformation must be invoked in order to avoid space problems. The models do illustrate that deformation (either localized or continuous and spatially distributed) to the south of the coseismic fault plane would produce uplift in a location to match the long-term uplift shown by the fluvial terraces.

Our study, along with the results of Copley (2014), Copley and Reynolds (2014), and Mackenzie et al. (2016), shows that it is not always valid to assume that fold growth above blind reverse faults is coseismic or that it is appropriate to model it as an elastic dislocation. Even in cases where folds and coseismic uplift have been seen to match topography, such as at the Coalinga anticline, Stein and Ekström (1992) reported that a modest fraction of the folding was postseismic, so at best modeling all fold growth as coseismic is likely to overestimate the coseismic moment release rate. One approach to extracting quantitative information from fold growth is to fit elastic dislocation models to terrace anticlines in an attempt to locate faults and estimate their slip rates (e.g., Benedetti et al., 2000). Our observational work, along with an earlier sandbox model by Bernard et al. (2007), suggests that this may not always be appropriate.

Many studies adopt a different approach to the analysis of folds above blind reverse faults by assuming a kinematic model and using that as the basis to derive slip rates (e.g., Amos et al., 2007; Daëron et al., 2007; Le Béon et al., 2014; Rockwell et al., 1988; Shaw et al., 2002). In subsequently estimating earthquake recurrence rates, the assumption is made that all the slip occurs in earthquakes, which as we have seen may not always be valid. The models are purely geometrical and do not need to specify physical mechanisms for the deformation, but the geometries implicitly require that not all of the fold growth occurs through coseismic slip in an elastic medium. This approach relies on a good choice of kinematic model which, as Hubbard et al. (2014) discuss, is not straightforward even at a location with a wealth of supplementary borehole and seismic reflection data. There is a need, therefore, to develop further physical models relating driving stresses to fold growth above blind faults.

8.2. Change in Fold Shape From Pliocene to Late Quaternary

The high-quality strike-perpendicular seismic reflection profile that directly crosses the coseismic slip patch allows us to study the long-term development of the fold and thrust structures at Pishan (section 4). The constant thickness of the Paleogene to early Pliocene sediments overlying the fault indicates that they predate the onset of folding at ≤ 3.6 Ma and therefore show the full cumulative displacement since that time (Jiang & Li, 2014; Li et al., 2016). They form a monocline with the southern limb dipping to the north slightly more steeply than the northern limb (Figure 6) and with a wavelength noticeably shorter than the topographic anticline—its two hinge axes are only 3 km apart (Figure 13). The anticline, though prominent in the surface geomorphology, is not visible in the long-term geological structure. This suggests that there was a change in shortening style before the formation and abandonment of the oldest terraces and recently enough for its effects to be below the detection level of the seismic reflection profile. The change in style is reminiscent of the northward shift of the topographic fold hinge from the older to younger terraces in Figure 8, and both stages of fold style development are likely to be related to the immaturity of the Hotan Fold and Thrust Belt. This highlights the importance of the balance of the shortening mechanisms behind the frontal thrust—the fault which broke in 2015 has accumulated significant offset and appears to have played an important role in the development of subsurface structure and to still play an important role in the present day, yet the style and location of folding above it have changed over time.

9. Conclusions

In summary, our seismic and geodetic results show that the Pishan earthquake occurred on the frontmost of a well-established stacked series (Li et al., 2016) of thrusts which does not reach the surface and had been independently mapped on seismic reflection data. The ability of InSAR to constrain the depth extent of faulting has proven to be useful affirmation of a velocity model in the depth migration of the seismic reflection survey and may be applicable elsewhere for the accurate imaging of thick sedimentary piles and structures with few or no prior velocity constraints. The first six months of postseismic deformation saw displacements of one fifth the magnitude of the coseismic static displacements, with the locations broadly similar but the spatial patterns differing slightly. We attribute the postseismic signal to afterslip below the coseismic slip patch and either afterslip at the top of the coseismic patch or distributed deformation in the unbroken overlying sediments.

The axis of the anticline in Quaternary terraces above the fault is offset by ~ 5 km back from the coseismic peak and ~ 6 km from the postseismic peak, confirming that the coseismic and early postseismic deformations are not sufficient to describe the growth of geological structures. Down to a depth of 6–8 km the topographic growth is accommodated by distributed deformation. We suggest that this deformation is influenced from below by aseismic deformation within or bounding the wedge or by the sum of further earthquakes around the 2015 slip patch. The difference in fold shape between the full cumulative record given by seismic reflection and Quaternary geomorphology shows that the style of shortening has changed recently in the lifetime of this range front.

By performing force balance calculations based upon earthquake depths and stress drops in and on the margins of the Tarim Basin, we have been able to estimate the strengths of the active faults and provide order-of-magnitude bounds on the viscosity of the ductile part of the Tibetan Plateau. The rheology contrast across the northern boundary of the Tibetan Plateau is likely to be responsible for the low dip angle of the Pishan fault plane.

Appendix A: InSAR Processing Methods

In interferometric wide swath mode, as used in this study, Sentinel-1 uses the Terrain Observation with Progressive Scans (TOPS) acquisition method (De Zan & Monti Guarnieri, 2006). This produces a swath with a total width of 250 km, acquired in three subswaths. Each subswath is acquired as a series of individual bursts of ~ 19 km length in which the beam is steered in the azimuth direction from backward to forward. Each full interferogram therefore contains multiple burst and subswath boundaries. Across each burst boundary there is a significant change in viewing geometry and Doppler centroid, so highly accurate coregistration is required to avoid phase jumps at the burst boundaries.

Here we used the GAMMA software to coregister single look complex (SLC) images to within <0.0005 SLC pixel in azimuth (less than 1 cm). An initial coregistration lookup table was made taking into account precise orbit state vectors and elevation using the 30 m SRTM1 DEM from the USGS. We then determined a uniform offset in range and azimuth to apply to the entire slave SLC. This was initially done using intensity cross correlation of a number of patches, iterated until the azimuth offset was <0.01 SLC pixel. We then refined the offset using a spectral diversity method (Scheiber & Moreira, 2000) and iterated until the azimuth offset was <0.0005 SLC pixel.

After coregistration we constructed the differential interferogram and filtered it 3 times using an adaptive spectral filter (Goldstein & Werner, 1998). This emphasizes the dominant wavelengths in the signal, preserving high-phase gradients and minimizing the extent of incoherent areas affecting unwrapping. The initial coherence was generally high, and interferograms were inspected visually after each step of filtering to ensure that there were no modifications during filtering which would bias our following analysis. We masked the interferograms to 95% coherence and unwrapped them using a Minimum Cost Flow algorithm. The final interferograms were geocoded and downsampled to 100 m. Where available we used the 30 m SRTM1 data set for topographic correction. We found that the voids in the SRTM data had been imperfectly filled in the 90 m SRTM Plus data set. At void boundaries there is a step change in elevation that led to phase errors in the interferograms. We therefore used the nonvoid filled SRTM1 DEM and did not retrieve the phase within the void areas. These regions are situated predominantly in the mountains to the south, well away from the tectonic deformation. The perpendicular baselines for all interferograms used in this study are <150 m.

Appendix B: Uniform-Slip Dislocation Model Method

To reduce computation times, and to give greater weight to the observations closer to the fault, we used a median filter to downsample our 100 m resolution interferograms to a 2 km grid in the near field and a 10 km grid in the far field. The interferograms do not contain any discontinuities as are seen in surface-rupturing earthquakes and, given the smooth deformation signal, a 2 km spacing gives high enough resolution to represent the full displacement signal given the spatial wavelength of correlated atmospheric noise. After downsampling we have 2,654 data points in total across the three tracks.

Our model comprises uniformly distributed slip on a rectangular dislocation in an elastic half-space. Elastic constants of the half-space were set by selecting Lamé parameters: $\lambda = \mu = 3.23 \times 10^{10}$ Pa, which are consistent with the velocities and density used in section 3.1. We used a nonlinear downhill Powell's algorithm with multiple Monte Carlo restarts (Wright et al., 1999) to find the best fitting model to our downsampled interferograms. Orbital errors and choice of unwrapping point were accounted for by solving for the nuisance parameters of linear gradients in phase across the interferograms (flattening) and a uniform phase shift.

We estimated uncertainties in our solution by making 250 perturbed interferograms which had noise added to them based on a characterization of the noise in a far field, minimally deformed, portion of each interferogram away from the earthquake (Funning et al., 2005; Wright et al., 2003). The fault parameters found by inverting these perturbed data sets are fitted with normal distributions (Figure S6) and give some indication of the influence of noise, mostly expected to be from the atmosphere (Scott & Lohman, 2016), on our results. This provides a more complete assessment of the uncertainties than those of He et al. (2016) and Wen et al. (2016). The standard deviation for each parameter is listed in Table 1.

Appendix C: Force Balance Methodology

The Pishan earthquake allows us to explore how the forces are transmitted between the Tarim Basin and the Tibetan Plateau. The forces between the two could be transmitted through a combination of faulting in the brittle crust (dark orange layer on Figure 12) and stresses exerted on the Tarim Basin lithosphere by deformation in the ductile middle and lower crust of Tibet (pale orange and yellow layers on Figure 12). The force transmitted across the range front thrust system, which broke in the Pishan earthquake, can be approximated as $F_b = W_f \tau$, where W_f is the downdip width of the faulting, and τ is the stress drop in the earthquake (we are here assuming that the stress drops are complete; if that is not the case, then the brittle layer could support a significantly higher proportion of the force). Copley and McKenzie (2007) argued that the surface motions in southern and eastern Tibet were consistent with a decrease in viscosity with depth, as would be expected for an increase in temperature with depth. We therefore calculate the stresses in the ductile part of the Tibetan crust assuming a two-layer viscosity structure (pale orange and yellow layers on Figure 12). This configuration makes the calculations simpler than using a viscosity structure that changes continuously with depth, but does not change our conclusions, our estimated viscosities for the upper (η_u) and lower (η_l) layer merely representing spatial averages. The force transmitted across the upper surface of the underthrusting Tarim lithosphere can be calculated as $F_{du} = W_d \eta_u V_t / H_u$, which is the consequence of the flow of Tibetan material over the Tarim Basin lithosphere, where W_d is the downdip width of Tarim lithosphere in contact with the ductile Tibetan lithosphere, V_t is the rate of convergence between northern Tibet and the Tarim Basin, and H_u is the vertical thickness of the ductile Tibetan lithosphere overlying the underthrust Tarim Basin lithosphere (Figure 12). The force on the "nose" of Tarim Basin can be estimated as $F_{dl} = H_l \eta_l V_i / D$, where H_l is the thickness of the underthrust Tarim Basin lithosphere, D is the distance between the rigid Indian and Tarim Basin lithospheres at depth, and V_i is the relative velocity between India and the Tarim Basin (Figure 12). F_{dl} represents the accommodation of the relative motion between the Tarim Basin and India at depth beneath the Tibetan Plateau. Given our estimate of the compressive stress above and that the forces must be balanced, $F_{total} = 4 \pm 2 \times 10^{12} \text{ Nm}^{-1} = F_b + F_{du} + F_{dl}$, we can therefore use our estimates of the geometry of the deformation, the stress drop in the Pishan earthquake, and GPS estimates of V_t and V_i to estimate the viscosities of the Tibetan upper and lower crust. Our calculation is completely instantaneous and represents the force balance that existed just before the Pishan earthquake occurred (i.e., we do not include the transient postseismic stress changes due to the event).

We estimated the stress drop of the Pishan earthquake to be 1–3 MPa by fitting a hemiellipsoid to our coseismic slip distribution and using the expressions of Bürgmann et al. (1994). We take V_t to be 2 mm/yr (Shen et al., 2001), V_i to be 15–20 mm/yr (Zhang et al., 2004), W_f to be 80 km (based on this study), and W_d to be 150 km

(Wittlinger et al., 2004). H_u is taken as 30 km and H_l as 40 km (Wittlinger et al., 2004). We have taken D to be 200 km; however, the true value could be lower (Craig et al., 2012). F_{dl} scales inversely with D , so errors in our estimate of D will result in errors in F_{dl} . Reducing D will have the effect of making F_{dl} more important compared with F_{du} . However, our aim is to provide order-of-magnitude estimates of the viscosities (which is all that is possible given our limited knowledge of the geometry and overall force balance), so given that errors in D scale inversely with F_{dl} , our results will not change dramatically for our largest likely errors in D . There is a trade-off between the values of η_u and η_l . If they are taken to be equal, then $\eta_u = \eta_l = 9 \pm 5 \times 10^{21}$ Pa s. If they are assumed to be different by 1 or 2 orders of magnitude, then either $\eta_u = 1.2 \pm 0.6 \times 10^{22}$ Pa s and $\eta_l = 1.2 \pm 0.6 \times 10^{21}$ Pa s or $\eta_u = 1.2 \pm 0.7 \times 10^{22}$ Pa s and $\eta_l = 1.2 \pm 0.7 \times 10^{20}$ Pa s. The final uncertainties on the viscosity quoted here incorporate uncertainties on the parameters as listed above.

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